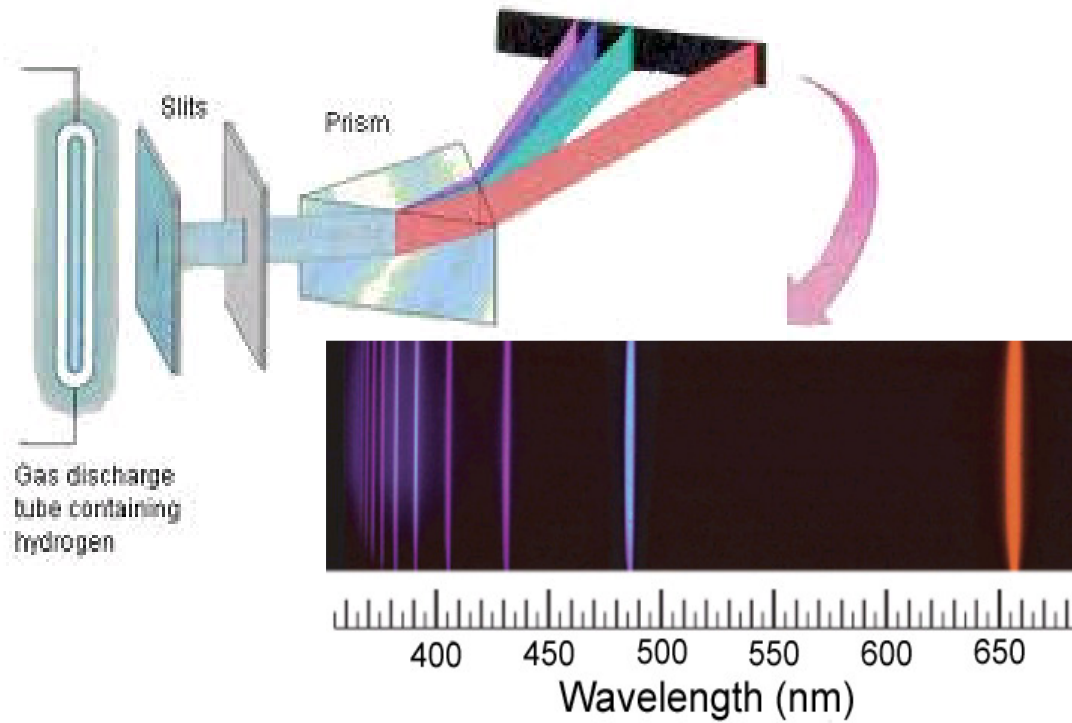


# Atomic spectrum of Hydrogen



Rydberg:

$$\frac{1}{\lambda} = R \left( \frac{1}{n_1^2} - \frac{1}{n_2^2} \right) \quad n_1, n_2 = 1, 2, 3, \dots$$

"Rydberg constant"  $\sim 10^5 \text{ cm}^{-1}$

$$\lambda \nu = c$$

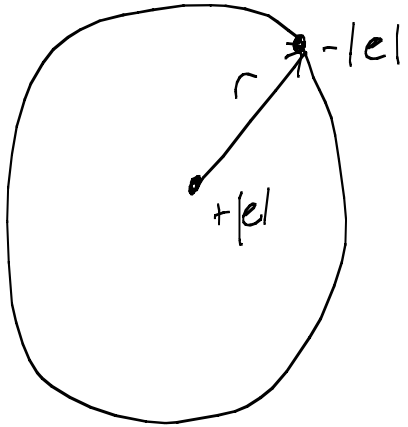
$$\lambda (h\nu) = hc$$

$$E = h\nu = \frac{hc}{\lambda} = \frac{1240 \text{ nm eV}}{\lambda}$$

$$R_y \rightarrow hcR \sim 13.6 \text{ eV}$$

"Rydberg"

## Classical Rutherford model



$$E_{\text{tot}} = E_{\text{kinetic}} + E_{\text{potential}} = \frac{1}{2}mv^2 + \left( -\frac{e^2}{4\pi\epsilon_0 r} \right) = \frac{1}{2}mv^2 - mv^2 = -\frac{1}{2}mv^2$$

$$F_{\text{centripetal}} = F_{\text{electrostatic}} \rightarrow \frac{mv^2}{r} = \frac{e^2}{4\pi\epsilon_0 r^2} \rightarrow r = \frac{e^2}{4\pi\epsilon_0 mv^2}$$

acceleration:  $\frac{v^2}{r} \rightarrow$  EM radiation and finite lifetime of atom!

# "Semiclassical" Bohr atom

de Broglie:  $\lambda = \frac{h}{p}$

Bohr: Circumference of orbit =  $2\pi r = n\lambda$   $n = 1, 2, 3 \dots$   
 $= n \frac{h}{mv}$

$$\rightarrow mvr = L = n \frac{h}{2\pi} = n\hbar$$

$$= mv \frac{e^2}{4\pi\epsilon_0 mv^2} = \frac{e^2}{4\pi\epsilon_0 v} = n\hbar \rightarrow v = \frac{e^2}{4\pi\epsilon_0 n\hbar}$$

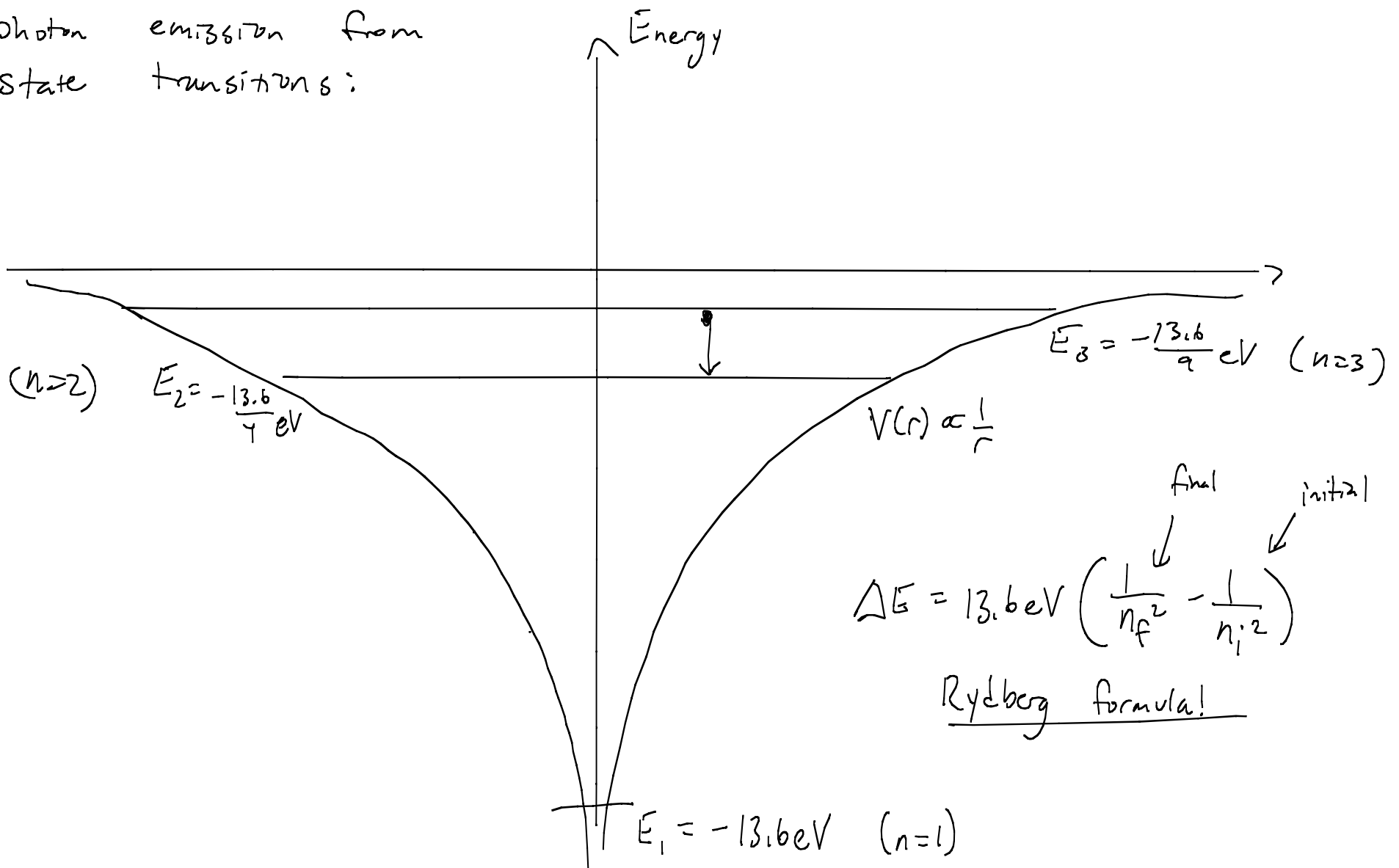
$$E_{\text{tot}} = -\frac{1}{2}mv^2 = -\frac{1}{2}m \left( \frac{e^2}{4\pi\epsilon_0 n\hbar} \right)^2 \approx -\frac{13.6 \text{ eV}}{n^2}$$

$$r = \frac{e^2}{4\pi\epsilon_0 mv^2} = \frac{4\pi\epsilon_0 \hbar^2 n^2}{me^2} = a_0 n^2$$

↑  
"Bohr radius"  $\sim 0.5 \text{ \AA}$  ( $5 \times 10^{-9} \text{ cm}$ )

# Hydrogen spectrum

photon emission from  
state transitions:



## Problems w/ Bohr's model of hydrogen

1. Orbit w/ definite radius violates uncertainty principle!

2. Angular momentum of ground state is non-zero!

3. Fails to explain:

a. spectra of larger atoms w/  $>1$  electron

b. multiplets (zero-field splitting of spectral lines)

c. magnetic field-induced spectral line splitting  
(Zeeman effect)

# Quantum mechanics of Hydrogen atom

$$V(r) = -\frac{e^2}{4\pi\epsilon_0 r} \quad (\text{spherically symmetric}) \quad \longrightarrow \quad \psi_{nlm}(r, \theta, \phi) = R_{nl}(r) Y_{lm}(\theta, \phi)$$

$$\text{Radial equation: } -\frac{\hbar^2}{2m} \frac{d^2 u}{dr^2} + \left[ -\frac{e^2}{4\pi\epsilon_0 r} + \frac{\hbar^2 l(l+1)}{2mr^2} \right] u = Eu, \quad u(r) \equiv rR(r)$$

$$k^2 \equiv -\frac{2mE}{\hbar^2}$$

$$\frac{1}{k^2} \frac{d^2 u}{dr^2} + \left[ \frac{e^2}{4\pi\epsilon_0 r \frac{\hbar^2 k^2}{2m}} - \frac{l(l+1)}{k^2 r^2} - 1 \right] u = 0$$

$$\rho \equiv kr \quad \rightarrow \quad d\rho = k dr$$

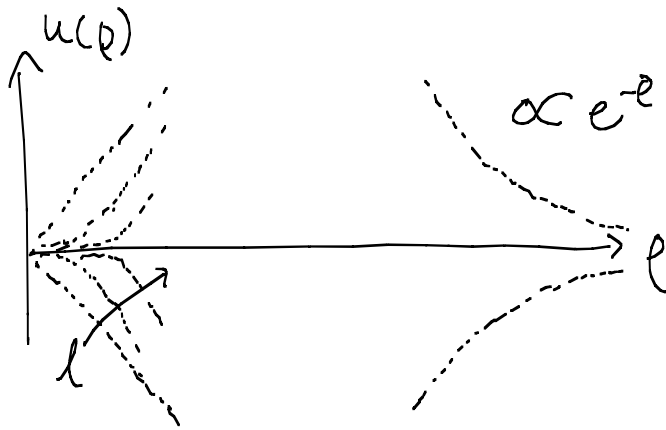
$$\frac{d^2 u}{d\rho^2} + \left[ \left( \frac{2me^2}{4\pi\epsilon_0 \hbar^2 k} \right) \frac{1}{\rho} - \frac{l(l+1)}{\rho^2} - 1 \right] u = 0$$

$$\frac{d^2 u}{d\rho^2} + \left[ \frac{\rho_0}{\rho} - \frac{l(l+1)}{\rho^2} - 1 \right] u = 0$$

## Asymptotic behavior of $u(\rho)$

$\rho \rightarrow \infty$ :  $\sim \frac{d^2 u}{d\rho^2} - u = 0 \rightarrow u(\rho) = A e^{-\rho} + \cancel{B e^{+\rho}}$  for normalizable wavefunction  
Boundary condition  $e \rho \rightarrow \infty$

$\rho \rightarrow 0$ :  $\sim \frac{d^2 u}{d\rho^2} - \frac{l(l+1)}{\rho^2} u = 0 \rightarrow u(\rho) = C \rho^{l+1} + \cancel{D \rho^{-l}}$  Boundary condition  $e \rho = 0$



Ansatz:  $u(\rho) = \rho^{l+1} e^{-\rho} v(\rho)$