

Solution To PHYS622 Homework Set 9

1 Solution1

$$\begin{aligned}
 V_{\pm}^{(1)} &= \mp \frac{V_x \pm iV_y}{\sqrt{2}}, V_0^{(1)} = V_z \\
 V_+^{(1)} &= -\frac{1}{\sqrt{2}}[(V_x \cos \beta + V_z \sin \beta) + iV_y] \\
 V_0^{(1)} &= -V_x \sin \beta + V_z \cos \beta \\
 V_{-1}^{(1)} &= \frac{1}{\sqrt{2}}[(V_x \cos \beta + V_z \sin \beta) - iV_y]
 \end{aligned}$$

If we rotate a vector $\vec{V}$ about the y-axis by an angle β , the new vector is:

$$\begin{bmatrix} V'_x \\ V'_y \\ V'_z \end{bmatrix} = \begin{bmatrix} \cos \beta & 0 & \sin \beta \\ 0 & 1 & 0 \\ -\sin \beta & 0 & \cos \beta \end{bmatrix} \begin{bmatrix} V_x \\ V_y \\ V_z \end{bmatrix}$$

2 Solution2

(a).

$$\begin{aligned}
 T_{+1}^{(1)} &= \frac{i}{2}[(U_y V_z - U_z V_y) + i(U_z V_x - V_z U_x)] \\
 T_0^{(1)} &= \frac{i}{\sqrt{2}i}[(U_x V_y - V_x U_y)] \\
 T_{-1}^{(1)} &= \frac{i}{2i}[(U_y V_z - U_z V_y) - i(U_z V_x - V_z U_x)]
 \end{aligned}$$

(b).

$$\begin{aligned}
 T_{+2}^{(2)} &= \frac{1}{2}[(U_x V_x - U_y V_y) + i(U_x V_y + V_x U_y)] \\
 T_{-2}^{(2)} &= \frac{1}{2}[(U_x V_x - U_y V_y) - i(U_x V_y + V_x U_y)] \\
 T_{+1}^{(2)} &= -\frac{1}{2}[(U_x V_z + U_z V_x) + i(U_y V_z + V_y U_z)] \\
 T_{-1}^{(2)} &= -\frac{1}{2}[(U_x V_z + U_z V_x) - i(U_y V_z + V_y U_z)] \\
 T_0^{(2)} &= \frac{1}{\sqrt{6}}[(2U_z V_z - U_x V_x - U_y V_y)]
 \end{aligned}$$

3 Solution3

$$A = \frac{l+1 + \vec{l} \cdot \vec{\sigma}}{2l+1}, B = \frac{l - \vec{l} \cdot \vec{\sigma}}{2l+1}$$

Note $\vec{l} \cdot \vec{\sigma} = \vec{J}^2 - \vec{l}^2 - \vec{S}^2$

$$A|j = l + 1/2, m \rangle = \frac{l + 1 + \vec{l} \cdot \vec{\sigma}}{2l + 1} |j = l + 1/2, m \rangle = \frac{l + 1 + \vec{J}^2 - \vec{l}^2 - \vec{S}^2}{2l + 1} |j = l + 1/2, m \rangle =$$

$$\frac{l + 1 + (l + 1/2)(l + 3/2) - l(l + 1) - 1/2(1/2 + 1)}{2l + 1} |j = l + 1/2, m \rangle = |j = l + 1/2, m \rangle$$

$$A|j = l - 1/2, m \rangle = 0$$

Similarly

$$B|j = l + 1/2, m \rangle = 0$$

$$B|j = l - 1/2, m \rangle = |j = l - 1/2, m \rangle$$

Write an arbitrary state in terms of the angular momentum states,

$$\psi = \sum_m C_m |j = l + 1/2, m \rangle + \sum_m D_m |j = l - 1/2, m \rangle$$

$$A\psi = \sum_m C_m |j = l + 1/2, m \rangle$$

$$A^2\Psi = A \sum_m C_m |j = l + 1/2, m \rangle = \sum_m C_m |j = l + 1/2, m \rangle$$

$$B\psi = \sum_m D_m |j = l - 1/2, m \rangle$$

$$B^2\Psi = B \sum_m C_m |j = l - 1/2, m \rangle = \sum_m C_m |j = l - 1/2, m \rangle$$

Hence, $A^2 = A$, $B^2 = B$.

$$\langle m - 1/2, 1/2 | l + 1/2, m \rangle = \sqrt{\frac{l + m + 1/2}{2l + 1}}$$

$$\langle m - 1/2, 1/2 | l - 1/2, m \rangle = \sqrt{\frac{l - m + 1/2}{2l + 1}}$$

$$\langle m + 1/2, -1/2 | l + 1/2, m \rangle = \sqrt{\frac{l - m + 1/2}{2l + 1}}$$

$$\langle m + 1/2, -1/2 | l - 1/2, m \rangle = \sqrt{\frac{l + m + 1/2}{2l + 1}}$$

4 Solution4

(a).

$$\langle n', l', m' | \mp \frac{1}{\sqrt{2}}(x \pm iy) | n, l, m \rangle = \langle l, 1; m, \pm 1 | l, 1; l', m' \rangle \frac{\langle n' l' || T^{(1)} || n l \rangle}{\sqrt{2l + 1}}$$

$$\begin{aligned}
& \langle n'l'm'|z|nlm \rangle = \langle l, 1; m, 0 | l, 1; l', m' \rangle \frac{\langle n'l' || T^{(1)} || nl \rangle}{\sqrt{2l+1}} \\
& \Rightarrow \frac{\langle n', l', m' | \mp \frac{1}{\sqrt{2}}(x \pm iy) | n, l, m \rangle}{\langle n'l'm'|z|nlm \rangle} = \frac{\langle l, 1; m, \pm 1 | l, 1; l', m' \rangle}{\langle l, 1; m, 0 | l, 1; l', m' \rangle}
\end{aligned}$$

(b).

$$\begin{aligned}
\langle n'l'm' | \mp \frac{1}{\sqrt{2}}(x \pm iy) | nlm \rangle &= \int \int \int R_{n'l'}^*(r) Y_{l'}^{*m'}(\theta, \phi) (\mp \frac{1}{\sqrt{2}})(x \pm iy) R_{nl}(r) Y_l^m(\theta, \phi) r^2 dr \sin \theta d\theta d\phi \\
\langle n'l'm' | z | nlm \rangle &= \int \int \int R_{n'l'}^*(r) Y_{l'}^{*m'}(\theta, \phi) (r \cos \theta) R_{nl}(r) Y_l^m(\theta, \phi) r^2 dr \sin \theta d\theta d\phi \\
&\Rightarrow \frac{\langle n'l'm' | \mp \frac{1}{\sqrt{2}}(x \pm iy) | nlm \rangle}{\langle n'l'm' | z | nlm \rangle} = \frac{\mp \frac{1}{\sqrt{2}} \int \int Y_{l'}^{*m'}(\theta, \phi) Y_l^m(\theta, \phi) e^{\pm i\phi} \sin \theta d\theta d\Omega}{\int \int Y_{l'}^{*m'}(\theta, \phi) Y_l^m(\theta, \phi) \cos \theta d\Omega}
\end{aligned}$$

Calculate the integration in the numerator and denominator, we have

$$\begin{aligned}
N &= \sqrt{\frac{2l+1}{2l'+1}} \langle l, 1; 0, 0 | l, 1; l', 0 \rangle \langle l, 1; m \pm 1 | l, 1; l', m' \rangle \\
D &= \sqrt{\frac{2l+1}{2l'+1}} \langle l, 1; 0, 0 | l, 1; l', 0 \rangle \langle l, 1; m, 0 | l, 1; l', m' \rangle
\end{aligned}$$

$$\frac{\langle n'l'm' | \mp \frac{1}{\sqrt{2}}(x \pm iy) | nlm \rangle}{\langle n'l'm' | z | nlm \rangle} = \frac{\langle l, 1; 0, 0 | l, 1; l', 0 \rangle \langle l, 1; m \pm 1 | l, 1; l', m' \rangle}{\langle l, 1; 0, 0 | l, 1; l', 0 \rangle \langle l, 1; m, 0 | l, 1; l', m' \rangle}$$

5 Solution5

(a).

$$\begin{aligned}
xy &= \frac{T_{+2}^{(2)} - T_{-2}^{(2)}}{2i} \\
xz &= \frac{T_{-1}^{(2)} - T_{+1}^{(2)}}{2} \\
x^2 - y^2 &= T_{+2}^{(2)} + T_{-2}^{(2)}
\end{aligned}$$

(b).

$$e \langle \alpha, j, m' | (x^2 - y^2) | \alpha, j, m = j \rangle = \frac{Q}{\sqrt{6}} \frac{\langle j, 2; m = j, q = -2 | j, 2; j, m' \rangle}{\langle j, 2; m = j, q = 0 | j, 2; j, m = j \rangle}$$

6 Solution6

$$H_{int} = A(3S_z^2 - S^2) + B(S_+^2 + S_-^2)$$

$$A = \frac{eQ}{4s(s-1)\hbar^2} \left(\frac{\partial^2 \phi}{\partial z^2} \right)_0$$

$$B = \frac{eQ}{4s(s-1)\hbar^2} \left[\left(\frac{\partial^2 \phi}{\partial x^2} \right)_0 + \frac{1}{2} \left(\frac{\partial^2 \phi}{\partial z^2} \right)_0 \right]$$

Write H_{int} as matrix form

$$H_{int} = \hbar^2 \begin{bmatrix} 3A & 0 & 2\sqrt{2} & 0 \\ 0 & -3A & 0 & 2\sqrt{3}B \\ 2\sqrt{3}B & 0 & -3A & 0 \\ 0 & 2\sqrt{3}B & 0 & 3A \end{bmatrix}$$

The eigenvalues are

$$E_{1,2} = -\sqrt{3}\sqrt{3A^2 + 4B^2}\hbar^2$$

$$E_{3,4} = \sqrt{3}\sqrt{3A^2 + 4B^2}\hbar^2$$

Two-fold degeneracy. Eigenvectors are

$$\begin{pmatrix} 0 \\ \frac{-3A + \sqrt{3}\sqrt{3A^2 + 4B^2}}{2\sqrt{3}B} \\ 0 \\ 1 \end{pmatrix}, \begin{pmatrix} \frac{-3A + \sqrt{3}\sqrt{3A^2 + 4B^2}}{2\sqrt{3}B} \\ 0 \\ 1 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ \frac{-3A + \sqrt{3}\sqrt{3A^2 + 4B^2}}{2\sqrt{3}B} \\ 0 \\ 1 \end{pmatrix}, \begin{pmatrix} \frac{3A + \sqrt{3}\sqrt{3A^2 + 4B^2}}{2\sqrt{3}B} \\ 0 \\ 1 \\ 0 \end{pmatrix}$$

7 Solution7

$$\vec{\mu} = \vec{\mu}_l + \vec{\mu}_s = \frac{e}{2mc}(\vec{L} + 2\vec{S})$$

$$\vec{\mu}^2 = \left(\frac{e}{2mc} \right)^2 [2\vec{J}^2 - \vec{L}^2 + 2\vec{S}^2]$$

$$\langle nljm_j | \vec{\mu}^2 | nljm_j \rangle = \left(\frac{e}{2m_j c} \right)^2 [2j(j+1) - l(l+1) + \frac{3}{2}] \hbar^2$$

$$\Rightarrow \mu = \left(\frac{e\hbar}{2m_j c} \right) \sqrt{2j(j+1) - l(l+1) + \frac{3}{2}}$$

From the projection theorem, we have

$$\langle nljm_j | \mu_0 | nljm_j \rangle = \frac{\langle nljm_j | \vec{J} \cdot \vec{\mu} | nljm_j \rangle}{j(j+1)\hbar^2} \langle jm_j | J_z | jm_j \rangle = \frac{m_j}{j(j+1)\hbar^2} \langle jljm_j | \vec{J} \cdot \vec{\mu} | nljm_j \rangle$$

$$\vec{J} \cdot \vec{\mu} = \frac{e}{2mc} [\vec{J}^2 + \vec{J} \cdot \vec{S}] = \frac{e}{2mc} \left[\frac{3}{2} \vec{J}^2 + \frac{1}{2} \vec{S}^2 - \frac{1}{2} \vec{L}^2 \right]$$

Therefore, we have

$$\langle nljm_j | \vec{J} \cdot \vec{\mu} | nljm_j \rangle = \frac{e}{2mc} \left[\frac{3}{2} j(j+1) - \frac{1}{2} l(l+1) + \frac{3}{8} \right] \hbar^2$$

$$\langle nljm_j | \mu_z | nljm_j \rangle = \frac{m_j}{j(j+1)} \frac{e}{2mc} \left[\frac{3}{2} j(j+1) - \frac{1}{2} l(l+1) + \frac{3}{8} \right]$$