

Solution To PHYS622 Homework Set 7

1 Solution1

$S_z(S_z + \hbar)(S_z - \hbar)$ must be vanish for spin 1 particle since $m = 0, 1, -1$ are the eigenvalue of S_z , and any spin state can be decomposed into the S_z eigenstates. Explicitly:

$$S_z = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} \hbar$$

$$S_z + \hbar = \begin{pmatrix} 2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \hbar$$

$$S_z - \hbar = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -2 \end{pmatrix} \hbar$$

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -2 \end{pmatrix} \hbar = 0$$

Similarly for

$$S_x = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \frac{\hbar}{\sqrt{2}},$$

we can rotate it to the S_x eigenstate and get exactly same result as for S_z .

2 Solution2

$$D^{1/2}(\alpha, \beta, \gamma) = \begin{pmatrix} e^{-i(\alpha+\gamma)/2} \cos \beta/2 & e^{-i(\alpha-\gamma)/2} \sin \beta/2 \\ e^{i(\alpha-\gamma)/2} \sin \beta/2 & e^{i(\alpha+\gamma)/2} \cos \beta/2 \end{pmatrix}.$$

Compare with (3.2.45)

$$e^{-\frac{\vec{\sigma} \cdot \vec{n}}{2} \theta} = \begin{pmatrix} \cos \theta/2 - i \hat{n}_z \sin \theta/2 & (-i n_x - n_y) \sin \theta/2 \\ (-i n_x + n_y) \sin \theta/2 & \cos \theta/2 + i \hat{n}_z \sin \theta/2 \end{pmatrix}$$

Therefore, we have

$$\begin{aligned} \cos \theta/2 &= \cos(\alpha + \gamma)/2 \cos \beta/2 \\ \sin \theta/2 &= \sqrt{1 - (\cos(\alpha + \gamma)/2 \cos \beta/2)^2} \\ n_z \sin \theta/2 &= \sin(\alpha + \gamma)/2 \cos \beta/2 \\ n_y \sin \theta/2 &= \cos(\alpha - \gamma)/2 \sin \beta/2 \\ n_x \sin \theta/2 &= -\sin(\alpha - \gamma)/2 \sin \beta/2 \end{aligned}$$

3 Solution3

(a).

$$J_+ J_- = (J_x + iJ_y)(J_x - iJ_y) = J_x^2 + J_y^2 + i(J_y J_x - J_x J_y) = \vec{J}^2 - J_z^2 + J_z \hbar$$

(b).

$$\begin{aligned} J_- \psi_{j,m} &= c_- \psi_{j,m-1} \\ J_+ \psi_{j,m}^* &= c_-^* \psi_{j,m-1}^* \\ \langle jm | J_+ J_- | jm \rangle &= |c_-|^2 = [j(j+1) - m^2 + m] \hbar^2 \end{aligned}$$

4 Solution4

$$\begin{aligned} L_{x,y} &= \frac{L_+ \pm L_-}{2} \\ \langle L_{x,y} \rangle &= \langle |L_+| \rangle \pm \langle |L_-| \rangle = 0 \end{aligned}$$

Since $L_{\pm}$ connect states with m to $m \pm 1$.

$$\begin{aligned} \langle lm | L_x^2 | lm \rangle &= \frac{1}{4} \langle lm | L_+^2 + L_-^2 + L_+ L_- + L_- L_+ | lm \rangle \\ &= \frac{1}{4} \langle lm | L_+ L_- + L_- L_+ | lm \rangle = \frac{1}{4} \langle lm | 2L_+ L_- - 2\hbar L_z | lm \rangle \\ &= \frac{\hbar^2}{2} [l(l+1) - m^2] \end{aligned}$$

Similarly, we can get $\langle L_y^2 \rangle = \langle L_x^2 \rangle$.

5 Solution5

$$D(R) = e^{-iJ_y \beta / \hbar}$$

The probability amplitude of the new state in $|l, m\rangle$ is

$$\begin{aligned} A_m &= \langle l, m | D(R) | 2, 0 \rangle = D_{m,0}^2(\beta) = \sqrt{\frac{4\pi}{2l+1}} Y_2^{m*}(\beta, 0) \\ A_0 &= \sqrt{\frac{4\pi}{5}} Y_2^0(\beta, 0) = \sqrt{\frac{4\pi}{5}} \sqrt{\frac{5}{16\pi}} (3 \cos^3 \beta - 1) \\ A_1 &= \sqrt{\frac{4\pi}{5}} Y_2^1(\beta, 0) = -\sqrt{\frac{4\pi}{5}} \sqrt{\frac{15}{8\pi}} \sin \beta \cos \beta \\ A_2 &= \sqrt{\frac{4\pi}{5}} Y_2^2(\beta, 0) = \sqrt{\frac{4\pi}{5}} \sqrt{\frac{15}{32\pi}} \sin^2 \beta \end{aligned}$$

So the probability:

$$\begin{aligned} |A_0|^2 &= \frac{1}{4}(3\cos^2\beta - 1)^2 \\ |A_1|^2 &= \frac{3}{2}(\sin\beta\cos\beta)^2 \\ |A_2|^2 &= \frac{3}{8}(\sin\beta)^4 \end{aligned}$$

6 Solution6

(a).

$$J_y = \frac{1}{2i}(J_+ - J_-)$$

$$\begin{aligned} &< j, m' | J_y | 1, m > \\ &= \frac{\hbar}{2i} [\sqrt{j(j+1) - m(m+1)} < j, m' | j, m+1 > - \sqrt{j(j+1) - m(m-1)} < j, m' | j, m-1 >] \end{aligned}$$

Therefore

$$J_y = \frac{\hbar}{2i} \left[\begin{pmatrix} 0 & \sqrt{2} & 0 \\ 0 & 0 & \sqrt{2} \\ 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 \\ \sqrt{2} & 0 & 0 \\ 0 & \sqrt{2} & 0 \end{pmatrix} \right] = \frac{i\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix}$$

(b).

$$e^{-iJ_y\beta/\hbar} = 1 - \frac{iJ_y\beta}{\hbar} + \frac{1}{2!} \left(\frac{iJ_y\beta}{\hbar} \right)^2 - \frac{1}{3!} \left(\frac{iJ_y\beta}{\hbar} \right)^3 + \dots$$

$$\begin{aligned} J_y^2 &= -\frac{\hbar^2}{2} \begin{pmatrix} -1 & 0 & 1 \\ 0 & -2 & 0 \\ 1 & 0 & -1 \end{pmatrix} \\ J_y^3 &= -\frac{i\hbar^2}{2\sqrt{2}} \begin{pmatrix} 0 & 2 & 0 \\ -2 & 0 & 2 \\ 0 & -2 & 0 \end{pmatrix} = \hbar^2 J_y \end{aligned}$$

So

$$J_y^4 = \hbar^2 J_y^2 \dots$$

$$\begin{aligned} e^{-iJ_y\beta/\hbar} &= 1 - \frac{i}{\hbar} J_y \left(\beta - \frac{\beta^3}{6} + \dots \right) - \frac{J_y^2}{\hbar^2} \left(\frac{\beta^2}{2!} - \frac{\beta^4}{4!} + \dots \right) \\ &= 1 - \frac{i}{\hbar} J_y \sin\beta - \left(\frac{J_y}{\hbar} \right)^2 (1 - \cos\beta) \end{aligned}$$

(c).

$$\begin{aligned}
d^{j=1}(\beta) &= e^{-iJ_y\beta/\hbar} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} - \frac{i}{\hbar} \sin\beta \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} \frac{i\hbar}{\sqrt{2}} + \frac{1 - \cos\beta}{2} \begin{pmatrix} -1 & 0 & 1 \\ 0 & -2 & 0 \\ 1 & 0 & -1 \end{pmatrix} \\
&= \begin{pmatrix} \frac{1+\cos\beta}{2} & \frac{-\sin\beta}{\sqrt{2}} & \frac{1-\cos\beta}{2} \\ \frac{\sin\beta}{\sqrt{2}} & \cos\beta & -\frac{\sin\beta}{\sqrt{2}} \\ \frac{1-\cos\beta}{2} & \frac{\sin\beta}{\sqrt{2}} & \frac{1+\cos\beta}{2} \end{pmatrix}
\end{aligned}$$