

Solution To PHYS622 Homework Set 5

1 Solution1

(a).

$$\begin{aligned}\psi(x) &= e^{if(x)/\hbar} \\ \hbar^2 \frac{\partial^2}{\partial x^2} \psi &= -p^2 \psi, \quad \frac{p^2}{2m} = E - V(x), \\ i\hbar f'' - f'^2 + p^2 &= 0.\end{aligned}$$

(b).

$$\begin{aligned}f(x) &= f_0(x) + \hbar f_1(x) + \hbar^2 f_2(x) + \dots \\ \text{therefore } f_0'^2 &= p^2, \quad if_0'' = 2f_0'f_1', \quad if_1'' = 2f_0'f_1' + f_1'^2\end{aligned}$$

(c).

$$\begin{aligned}f_0 &= \int^x \sqrt{2m(E - V(x))} dx \\ f_1' &= if_0''/2f_0' \\ \text{Therefore } f_1 &= i/2 \ln(f_0') = i/4 \ln(2m(E - V(x))), \\ \psi &\sim \frac{1}{[E - V(x)]^2} e^{\pm \int^x \sqrt{2m(E - V(x))} dx}\end{aligned}$$

2 Solution2

For a harmonic oscillator

$$V(x) = 1/2m\omega^2 x^2$$

The WKB solution within the potential will should be

$$\begin{aligned}2C/\sqrt{p} &= \cos(-1/\hbar \int_x^{x_2} pdx + \pi/4) \\ 2D/\sqrt{p} &= \cos(1/\hbar \int_x^{x_2} pdx - \pi/4)\end{aligned}$$

These two solutions are consistent if C=D, and

$$\begin{aligned}-1/\hbar \int_x^{x_2} pdx - \pi/4 &= 1/\hbar \int_x^{x_2} pdx - \pi/4 + n\pi \\ \text{therefore } \int_{x_1}^{x_2} \sqrt{2m(E - V(x))} dx &= \hbar(n + 1/2)\pi, \quad n = 0, 1, \dots\end{aligned}$$

$$\text{Let } x = \sqrt{\frac{2E}{m\omega^2}} \cos\theta, \quad \int_0^\pi \frac{2E}{\omega} \sin^2\theta d\theta = (n + \frac{1}{2})\hbar\pi \rightarrow \frac{2E}{\omega} \frac{\pi}{2} = (n + \frac{1}{2})\hbar\pi \rightarrow E_n = (n + \frac{1}{2})\hbar\omega.$$

3 Solution3

For a finite square barrier of height $V_0 > E$ and width $2a$, the WKB solution within the barrier is

$$\psi_{WKB} = \frac{A}{\sqrt{|p|}} e^{\frac{1}{\hbar} \int_0^x |p(x')| dx'} + \frac{B}{\sqrt{|p|}} e^{-\frac{1}{\hbar} \int_0^x |p(x')| dx'}$$

For approximation we can let $A = 0$, the solution near $x = 2a$ is

$$\begin{aligned}\psi(2a) &\sim \frac{1}{\sqrt{|p|}} e^{-\frac{1}{\hbar} \int_0^{2a} \sqrt{2m(V_0 - E)}} \\ T &= \frac{|\psi(2a)|^2}{|\psi(0)|^2} \sim e^{-\frac{4a}{\hbar} \sqrt{2m(V_0 - E)}}\end{aligned}$$

4 Solution4

$$-\frac{\hbar^2}{2m} \psi'' - V_0 \delta(x) \psi = -E \psi$$

There is one unique bound state

$$\psi(x) = A e^{\pm \sqrt{\frac{2mE}{\hbar^2}} x}, x > 0, x < 0.$$

Because we have two boundary conditions:

$$\begin{aligned}\psi(x)|_{x=0+} &= \psi(x)|_{x=0-} \\ \psi(x)|_{x=0+} - \psi(x)|_{x=0-} + \frac{2mV_0}{\hbar^2} \psi(0) &= 0 \\ \text{Therefore } E &= V_0^2 m / 2\hbar^2, A = \sqrt{mV_0/\hbar^2}.\end{aligned}$$

If the potential is switched off at $t = 0$, the wavefunction for $t > 0$ will be

$$\psi(x, t > 0) = \int_{-\infty}^{+\infty} dx' \psi(x', 0) K(x, x'; t, 0) = \sqrt{mV_0/\hbar^2} \sqrt{m/2\pi i \hbar} \int_{-\infty}^{+\infty} dx' e^{-mV_0/\hbar^2 + i(x-x')^2 m/2\hbar t}$$

5 Solution5

$$K(\vec{p}'', \vec{p}'; t, t_0) = \sum_{a'} \langle \vec{p}'' | a' \rangle \langle a' | \vec{p}' \rangle e^{-\frac{iE(a')(t-t_0)}{\hbar}}$$

Take a' to be p

$$K(\vec{p}'', \vec{p}'; t, t_0) = \sum_p e^{-\frac{iE(p')(t-t_0)}{\hbar}} \delta(\vec{p}'' - \vec{p}) \delta(\vec{p} - \vec{p}') = e^{-\frac{p'^2(t-t_0)}{2m\hbar}} \delta(\vec{p}'' - \vec{p}')$$

6 Solution6

Classical action

$$S = \int dt \left[\frac{m\dot{x}}{2} - V(x) \right]$$

$$S(n, n-1) = \Delta t \left[\frac{m}{2} \left(\frac{(x_n - x_{n-1})^2}{\Delta t} - \frac{\omega^2(x_n + x_{n-1})^2}{8} \right) \right]$$

$$\langle x_n, t_n | x_{n-1}, t_{n-1} \rangle = \sqrt{\frac{m}{2\pi\hbar\Delta t}} e^{\frac{iS(n,n-1)}{\hbar}}$$

The propagator is

$$K(x'', t; x', t_0) = \sqrt{\frac{m\omega}{2\pi i\hbar \sin(\omega(t-t_0))}} e^{\frac{im\omega[(x''^2+x'^2)\sin\omega(t-t_0)-2x'x'']}{2\hbar\sin\omega(t-t_0)}}$$

In the limit $\Delta t \rightarrow 0$

$$\sin\omega\Delta t \sim \omega\Delta t, \cos\omega\Delta t \sim 1 - \omega^2\Delta t^2/2$$

$$\text{therefore } K \rightarrow \sqrt{\frac{m\omega}{2\pi i\hbar\Delta t}} e^{\frac{im}{2\hbar\Delta t} [(x_n - x_{n-1})^2 - \frac{2x_n^2\omega^2\Delta t^2}{2}]}$$

$$\sim \sqrt{\frac{m\omega}{2\pi i\hbar\Delta t}} e^{\frac{im}{2\hbar\Delta t} [(x_n - x_{n-1})^2 - \frac{\omega^2(x_n + x_{n-1})^2}{8}]}$$