

Lecture 31 Highlights

Quantum statistical mechanics, continued...

We found the statistical weight W of the arrangement $(n_1, n_2, n_3, n_4, \dots, n_s, \dots)$:

$$W(n_1, n_2, \dots, n_s, \dots) = \prod_{s=1}^{\infty} P_s,$$

for three different types of statistics. This weight is proportional to the probability of finding this particular distribution of occupation numbers.

1) Distinguishable classical particles $W_{Dist}(n_1, n_2, \dots, n_s, \dots) = N! \prod_{s=1}^{\infty} \frac{g_s^{n_s}}{n_s!}$ (1)

2) Indistinguishable identical Fermions

$$W_{Fermions}(n_1, n_2, \dots, n_s, \dots) = \prod_{s=1}^{\infty} \frac{g_s!}{n_s!(g_s - n_s)!} \quad (2)$$

3) Indistinguishable identical Bosons $W_{Bosons}(n_1, n_2, \dots, n_s, \dots) = \prod_{s=1}^{\infty} \frac{(n_s + g_s - 1)!}{n_s!(g_s - 1)!}$ (3)

The next step is to maximize $W(n_1, n_2, \dots, n_s, \dots)$ by varying all of the occupation numbers, subject to the number and total energy constraints: $\sum_{i=1}^{\infty} n_i = N$ and $\sum_{i=1}^{\infty} n_i E_i = E$.

We will include the constraints using the method of Lagrange multipliers. This method allows one to perform a constrained maximization. We will form a new function to maximize, namely;

$$G(n_1, n_2, \dots, n_s, \dots, \alpha, \beta) = W(n_1, n_2, \dots, n_s, \dots) + \alpha \left(N - \sum_{i=1}^{\infty} n_i \right) + \beta \left(E - \sum_{i=1}^{\infty} n_i E_i \right)$$

To maximize this function we must enforce these conditions:

$$\frac{\partial G}{\partial n_s} = 0 \quad \forall s \quad \text{and} \quad \frac{\partial G}{\partial \alpha} = \frac{\partial G}{\partial \beta} = 0.$$

The form of G already satisfies the last two conditions.

Because of the products appearing in Eqs. (1)-(3), it is easier to maximize the logarithm of W . This will yield the same result since W and $\ln W$ have maxima at the same values of their arguments. The newly defined G for distinguishable particles now is:

$$\begin{aligned} G_{Dist}(n_1, n_2, \dots, n_s, \dots, \alpha, \beta) &= \ln \left(N! \prod_{s=1}^{\infty} \frac{g_s^{n_s}}{n_s!} \right) + \alpha \left(N - \sum_{i=1}^{\infty} n_i \right) + \beta \left(E - \sum_{i=1}^{\infty} n_i E_i \right) \\ &= \ln N! + \sum_{s=1}^{\infty} (n_s \ln g_s - \ln n_s!) + \alpha \left(N - \sum_{i=1}^{\infty} n_i \right) + \beta \left(E - \sum_{i=1}^{\infty} n_i E_i \right) \end{aligned}$$

To take the derivative of G with respect to n_s we must now decide what to do with the logarithm of $n_s!$. One approach is to employ Stirling's approximation:

$\ln x! \cong x \ln x - x$, good for $x \gg 1$. With this approximation, G_{Dist} becomes:

$$G_{Dist}(n_1, n_2, \dots, n_s, \dots, \alpha, \beta) \cong \ln N! + \sum_{s=1}^{\infty} (n_s \ln g_s - n_s \ln n_s + n_s) + \alpha \left(N - \sum_{i=1}^{\infty} n_i \right) + \beta \left(E - \sum_{i=1}^{\infty} n_i E_i \right)$$

Taking the derivative of G with respect to some particular n_s (called n_i in the lecture) and setting it equal to zero (to find the maximum), yields;

$$n_s = g_s e^{-(\alpha + \beta E_s)} \quad \text{Distinguishable particles}$$

For the other cases one gets

$$n_s = \frac{g_s}{e^{+(\alpha + \beta E_s)} + 1} \quad \text{Identical Fermions}$$

$$n_s = \frac{g_s}{e^{+(\alpha + \beta E_s)} - 1} \quad \text{Identical Bosons}$$

What are the Lagrange multipliers α and β ? They are determined by the number and energy constraints $\sum_{i=1}^{\infty} n_i = N$ and $\sum_{i=1}^{\infty} n_i E_i = E$. . The challenge is to determine the

energies and degeneracies of all of the single-particle states of the system - this is the hardest part of quantum statistical mechanics. Calculating the total energy of an ideal gas, which is a relatively easy case, Griffiths (pp. 239-240) finds that $\beta = 1/k_B T$, where T is the absolute temperature of the gas. The other parameter α is re-defined in terms of the chemical potential μ as $\alpha \equiv -\mu\beta$. The chemical potential is a measure of how much energy is required to change the particle number of the system from N to $N + 1$. The three distribution functions can now be written as:

$$n_s = g_s e^{-(E_s - \mu)/k_B T} \quad \text{Distinguishable particles}$$

$$n_s = \frac{g_s}{e^{+(E_s - \mu)/k_B T} + 1} \quad \text{Identical Fermions}$$

$$n_s = \frac{g_s}{e^{+(E_s - \mu)/k_B T} - 1} \quad \text{Identical Bosons}$$