

Fourier's Trick relies upon the following mathematical identity:

$$\int_0^L \sin\left(\frac{m\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) dx = \begin{cases} 0, & m \neq n \\ \frac{L}{2}, & m = n \end{cases}$$

We like to write this more compactly:

Define $\delta_{mn} \equiv \begin{cases} 0, & m \neq n \\ 1, & m = n \end{cases}$ "Kronecker Delta"

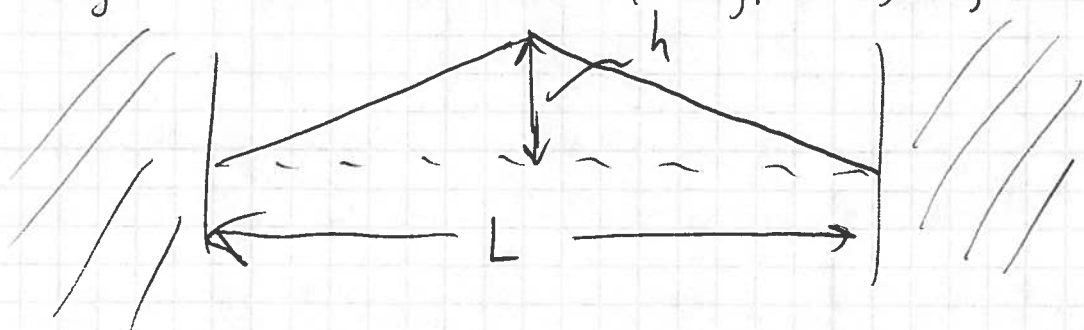
Then $\delta_{11} = 1$, $\delta_{12} = 0$, $\delta_{13} = 0$, $\delta_{22} = 1$, etc

Using the Kronecker Delta we can say

$$\int_0^L \sin\left(\frac{m\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) dx = \frac{L}{2} \delta_{mn}$$

or $\frac{2}{L} \int_0^L \sin\left(\frac{m\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) dx = \delta_{mn}$

Let's go back to our triangular string:



This is the shape at $t = 0$. The functional form is

$$y(x, t=0) = \begin{cases} \left(\frac{zh}{L}\right)x & , 0 \leq x \leq \frac{L}{2} \\ \left(\frac{zh}{L}\right)(L-x) & , \frac{L}{2} \leq x \leq L \end{cases}$$

We want to describe this simple function in a much more complicated way: as an infinite sum of normal modes:

$$y(x, t=0) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right)$$

The question is: what are the $\{a_n\}$?

Fourier's Trick tells us that any particular coefficient, ~~can be calculated~~ for example, the m^{th} coefficient (a_m), can be calculated by evaluating this integral:

$$a_m = \frac{2}{L} \int_0^L \sin\left(\frac{m\pi x}{L}\right) y(x) dx$$

For our function $y(x)$, this integral is

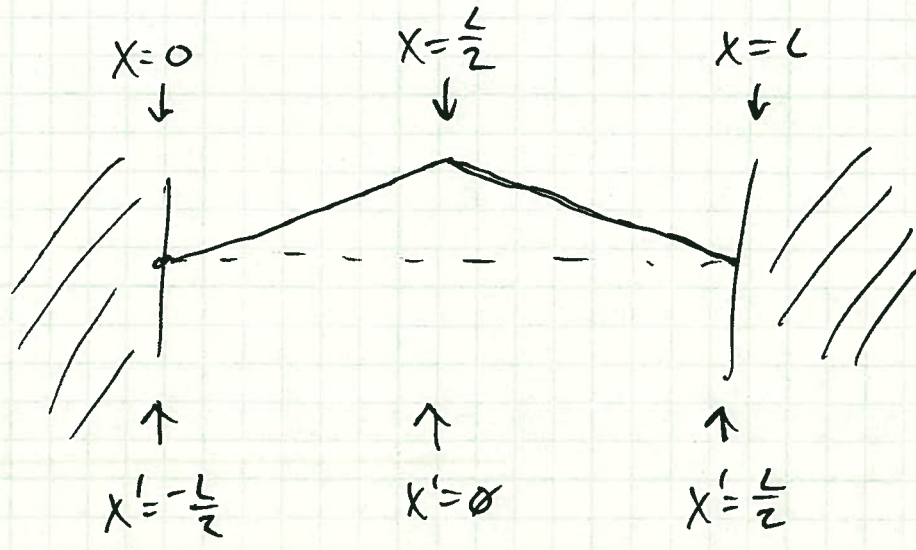
$$a_m = \frac{2}{L} \int_0^{L/2} \sin\left(\frac{m\pi x}{L}\right) \left(\frac{zhx}{L}\right) dx + \frac{2}{L} \int_{L/2}^L \sin\left(\frac{m\pi x}{L}\right) \left(\frac{zh(L-x)}{L}\right) dx$$

It turns out that the easiest way to evaluate this integral is to move our coordinate system -

$$\text{Let } x' \equiv x - \frac{L}{2}$$

$$\text{so that } x = x' + \frac{L}{2}$$

This means that $x' = 0$ is the center of the string



In terms of x' , our string position at $t=0$ is

$$y(x', t=0) = \begin{cases} \left(\frac{2h}{L}\right)\left(x' + \frac{L}{2}\right) & , \quad -\frac{L}{2} \leq x' \leq 0 \\ \left(\frac{2h}{L}\right)\left(-x' + \frac{L}{2}\right) & , \quad 0 \leq x' \leq \frac{L}{2} \end{cases}$$

Note that y is an even function of x' .

Also, we have the following math theorem:

IF $x = x' + \frac{L}{2}$,

Then $\sin\left(\frac{m\pi x}{L}\right) = \begin{cases} (-1)^{(m-1)/2} \cos\left(\frac{m\pi x'}{L}\right) & , \text{ for } m = \text{odd} \\ (-1)^{m/2} \sin\left(\frac{m\pi x'}{L}\right) & , \text{ for } m = \text{even} \end{cases}$

Now our integral has 2 cases:

$$a_m = \frac{2}{L} \int_{-\frac{L}{2}}^{\frac{L}{2}} y(x') (-1)^{(m-1)/2} \cos\left(\frac{m\pi x'}{L}\right) dx' \quad \text{for } m = \text{odd}$$

AND

$$a_m = \frac{2}{L} \int_{-\frac{L}{2}}^{\frac{L}{2}} y(x') (-1)^{(m)/2} \sin\left(\frac{m\pi x'}{L}\right) dx' \quad \text{for } m = \text{even}$$

This integrand is an ~~even~~ odd function of x' , because $y(x')$ is even, and $\sin\left(\frac{m\pi x'}{L}\right)$ is odd.

Therefore the integral is zero because we integrate from $-\frac{L}{2}$ to $\frac{L}{2}$.

So we only need to evaluate the case for $m = \text{odd}$:

$$a_m = \frac{2}{L} \int_{-\frac{L}{2}}^{\frac{L}{2}} y(x') (-1)^{(m-1)/2} \cos\left(\frac{m\pi x'}{L}\right) dx' \quad , \quad m = \text{odd.}$$

This integrand is even because $y(x')$ and $\cos\left(\frac{m\pi x'}{L}\right)$ are both even functions of x' . Since we integrate from $-\frac{L}{2}$ to $\frac{L}{2}$, we can just integrate from zero to $\frac{L}{2}$ and multiply by 2:

$$a_m = (2) \frac{2}{L} \int_0^{\frac{L}{2}} y(x') (-1)^{(m-1)/2} \cos\left(\frac{m\pi x'}{L}\right) dx'$$

$$a_m = (2) \left(\frac{2}{L}\right) (-1)^{(m-1)/2} \left(\frac{2h}{L}\right) \int_0^{L/2} \left(-x' + \frac{L}{2}\right) \cos\left(\frac{m\pi x'}{L}\right) dx'$$

$$= \left(\frac{8h}{L^2}\right) (-1)^{(m-1)/2} \left[\left(-\left(\frac{L}{m\pi}\right)^2 \cos\left(\frac{m\pi x'}{L}\right) - \frac{x' L}{m\pi} \sin\left(\frac{m\pi x'}{L}\right) + \left(\frac{L}{2}\right) \left(\frac{L}{m\pi}\right) \sin\left(\frac{m\pi x'}{L}\right) \right) \Big|_0^{L/2} \right]$$

zero for m=odd cancel

$$= \left(\frac{8h}{L^2}\right) (-1)^{(m-1)/2} \left[-\left(\frac{L}{m\pi}\right)^2 \cos\left(\frac{m\pi}{2}\right) - \frac{L^2}{2m\pi} \sin\left(\frac{m\pi}{2}\right) + \left(\frac{L^2}{2m\pi}\right) \sin\left(\frac{m\pi}{2}\right) \right]$$

* - $\left(-\left(\frac{L}{m\pi}\right)^2\right)$

$a_m = \frac{8h}{(m\pi)^2} (-1)^{(m-1)/2}$	$m = \text{odd}$	and	$a_m = \emptyset$ for m-even
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Therefore $a_1 = \frac{8h}{\pi^2}$

$a_2 = \emptyset$

$a_3 = -\frac{8h}{9\pi^2}$

$a_4 = \emptyset$

$a_5 = \frac{8h}{25\pi^2}$

⋮

Or we can write

$$y(x, t=0) = \frac{8h}{\pi^2} \sin\left(\frac{\pi x}{L}\right) - \frac{8h}{9\pi^2} \sin\left(\frac{3\pi x}{L}\right) + \frac{8h}{25\pi^2} \sin\left(\frac{5\pi x}{L}\right) + \dots$$

Why did we do this?

Recall our motivation: The general solution to the wave equation is a sum over normal modes:

$$\frac{\partial^2 y}{\partial x^2} = \frac{\partial^2 y}{\partial t^2} \Rightarrow y(x, t) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{L}\right) e^{i\omega_n t}$$

where $c_n = \text{real and imaginary}$
 $\equiv a_n + i b_n$

Given the initial condition:

$$y(x, t=0) = \text{a triangle} = \begin{cases} \left(\frac{2h}{L}\right)x, & 0 \leq x \leq \frac{L}{2} \\ \frac{2h}{L}(L-x), & \frac{L}{2} \leq x \leq L \end{cases}$$

We found the a_n :

$$a_n = \begin{cases} \left(\frac{8h}{(n\pi)^2}\right) (-1)^{(n-1)/2} & \text{for } n = \text{odd} \\ \emptyset & \text{for } n = \text{even} \end{cases}$$

What about the imaginary part, $\{b_n\}$?

It is determined by the initial velocity:

$$y(x, t=0) = \sum_{n=1}^{\infty} -\omega_n b_n \sin\left(\frac{n\pi x}{L}\right)$$

If we release the string from rest, then

$$\text{we must have } y(x, t=0) = 0 \Rightarrow \boxed{b_n = 0 \text{ for all } n}$$

So our final, time-dependent solution is

$$y(x, t) = \sum_{\substack{n=1 \\ \text{(only odd } n)}}^{\infty} \frac{8h}{(n\pi)^2} (-1)^{\frac{n-1}{2}} \sin\left(\frac{n\pi x}{L}\right) e^{i\omega_n t}$$

(odd n only)

$$\text{where } \omega_n = \sqrt{\frac{T}{\rho}} \frac{n\pi}{L}, \quad n=1, 2, 3, \dots$$

We write the initial condition function $y(x, t=0)$ as a sum over normal modes because then the time development is extremely simple: each normal mode goes forward in time with its own harmonic factor ($e^{i\omega_n t}$).