

PHYS 373 (Fall 2015): Mathematical Methods for Physics II

Summary of topics/formulae for Final exam

Chapter 7 of Boas (Fourier Series and Transforms)

1. General Fourier series: a function $f(x)$ with period $2l$ can be expanded as

$$f(x) = \frac{a_0}{2} + \sum_1^{\infty} \left(a_n \cos \frac{n\pi x}{l} + b_n \sin \frac{n\pi x}{l} \right) \quad (1)$$

$$= \sum_{-\infty}^{\infty} c_n e^{in\pi x/l} \quad (2)$$

where the Fourier coefficients are given by

$$a_n = \frac{1}{l} \int_{-l}^l f(x) \cos \frac{n\pi x}{l} dx \quad (3)$$

$$b_n = \frac{1}{l} \int_{-l}^l f(x) \sin \frac{n\pi x}{l} dx \quad (4)$$

$$c_n = \frac{1}{2l} \int_{-l}^l f(x) e^{-in\pi x/l} dx \quad (5)$$

2. Special Fourier series: we have

$$\text{if } f(x) \text{ is odd, } \begin{cases} b_n = \frac{2}{l} \int_0^{\infty} f(x) \sin \frac{n\pi x}{l} dx \\ a_n = 0 \end{cases} \quad (6)$$

$$\text{if } f(x) \text{ is even, } \begin{cases} a_n = \frac{2}{l} \int_0^{\infty} f(x) \cos \frac{n\pi x}{l} dx \\ b_n = 0 \end{cases} \quad (7)$$

3. Parseval's theorem for Fourier series:

$$\text{The average of } |f(x)|^2 \text{ (over a period)} = \left(\frac{1}{2} a_0 \right)^2 + \frac{1}{2} \sum_1^{\infty} a_n^2 + \frac{1}{2} \sum_1^{\infty} b_n^2 \quad (8)$$

$$= \sum_{-\infty}^{\infty} |c_n|^2 \quad (9)$$

4. General Fourier transform

$$f(x) = \int_{-\infty}^{\infty} g(\alpha) e^{i\alpha x} d\alpha \quad (10)$$

where

$$g(\alpha) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) e^{-i\alpha x} dx \quad (11)$$

5. Special Fourier transform: for an odd function, we have

$$f_s(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} g_s(\alpha) \sin \alpha x \, d\alpha \quad (12)$$

$$g_s(\alpha) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f_s(x) \sin \alpha x \, dx \quad (13)$$

Similarly, for an even function:

$$f_c(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} g_c(\alpha) \cos \alpha x \, d\alpha \quad (14)$$

$$g_c(\alpha) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f_c(x) \cos \alpha x \, dx \quad (15)$$

$$(16)$$

6. Parseval's theorem for Fourier transform:

$$\int_{-\infty}^{\infty} |g(\alpha)|^2 d\alpha = \int_{-\infty}^{\infty} \frac{1}{2\pi} |f(x)|^2 dx \quad (17)$$

Chapter 3 of Boas (Linear Algebra)

1. n -dimensional vector-space:

$$\mathbf{A} \cdot \mathbf{B} \text{ (inner product)} = \sum_1^n A_i B_i \quad (18)$$

$$A \text{ (norm)} = \sqrt{\mathbf{A} \cdot \mathbf{A}} \quad (19)$$

$$\mathbf{A} \text{ and } \mathbf{B} \text{ are orthogonal if } \mathbf{A} \cdot \mathbf{B} = 0 \quad (20)$$

2. vector-space of functions on $a \leq x \leq b$:

$$\text{Inner product of } A(x) \text{ and } B(x) = \int_a^b A^*(x) B(x) dx \quad (21)$$

$$\text{Norm of } A(x) = \sqrt{\int_a^b A^*(x) A(x) dx} \quad (22)$$

$$A(x) \text{ and } B(x) \text{ are orthogonal if } \int_a^b A^*(x) B(x) dx = 0 \quad (23)$$

3. Gram-Schmidt method for making a basis (\mathbf{A} , \mathbf{B} , \mathbf{C} ...) orthonormal:

$$\mathbf{e}_1 = \frac{\mathbf{A}}{A} \quad (24)$$

$$\mathbf{e}_2 = \text{normalized } [\mathbf{B} - (\mathbf{e}_1 \cdot \mathbf{B}) \mathbf{e}_1] \quad (25)$$

$$\mathbf{e}_3 = \text{normalized } [\mathbf{C} - (\mathbf{e}_1 \cdot \mathbf{C}) \mathbf{e}_1 - (\mathbf{e}_2 \cdot \mathbf{C}) \mathbf{e}_2] \quad (26)$$

Chapter 12 of Boas (Series Solutions of Differential Equations)

1. Series method for solving (linear) ordinary differential equations (ODE): assume a solution of the form (with a 's being coefficients *to be found*)

$$y = \sum_{n=0}^{\infty} a_n x^n \quad (27)$$

giving

$$y' = \sum_{n=0}^{\infty} n a_n x^{n-1} \quad (28)$$

$$y'' = \sum_{n=0}^{\infty} n(n-1) a_n x^{n-2} \quad (29)$$

Plug the above series into each term of the ODE. Find the *total* coefficient of each power of x on each side of ODE and equate them (again, for each power of x). This will give the higher a coefficients in terms of lower ones.

2. Legendre's equation:

$$(1-x^2)y'' - 2xy' + l(l+1)y = 0 \quad (30)$$

has a solutions for each integer l (chosen to be non-negative) which is called the Legendre polynomial, $P_L(x)$ defined with

$$P_l(1) = 1 \quad (31)$$

For example,

$$P_0(x) = 1, \quad P_1(x) = x, \quad P_2(x) = \frac{1}{2}(3x^2 - 1), \quad P_3(x) = \frac{1}{2}(5x^3 - 3x) \dots \quad (32)$$

3. Rodrigues' formula for Legendre polynomials

$$P_l(x) = \frac{1}{2^l l!} \frac{d^l}{dx^l} (x^2 - 1)^l \quad (33)$$

4. Generating function for Legendre polynomials:

$$\Phi(x, h) = (1 - 2xh + h^2)^{-1/2}, \quad |h| < 1 \quad (34)$$

$$= \sum_{l=0}^{\infty} h^l P_l(x) \quad (35)$$

5. Recursion relations for Legendre polynomials:

$$lP_l(x) = (2l-1)xP_{l-1}(x) - (l-1)P_{l-2}(x), \quad (36)$$

$$xP_l'(x) - P_{l-1}'(x) = lP_l(x), \quad (37)$$

$$P_l'(x) - xP_{l-1}'(x) = lP_{l-1}(x), \quad (38)$$

$$(1-x^2)P_l'(x) = lP_{l-1}(x) - lxP_l(x), \quad (39)$$

$$(2l+1)P_l(x) = P_{l+1}'(x) - P_{l-1}'(x), \quad (40)$$

$$(1-x^2)P_{l-1}'(x) = lxP_{l-1}(x) - lP_l(x) \quad (41)$$

6. Orthogonality of Legendre polynomials:

$$\int_{-1}^1 P_l(x)P_m(x)dx = 0, \text{ unless } l = m \quad (42)$$

7. Normalization of Legendre polynomials:

$$\int_{-1}^1 [P_l(x)]^2 = \frac{2}{2l+1} \quad (43)$$

$$(44)$$

8. A function defined over the interval $(-1, 1)$ can be expanded in a Legendre series

$$f(x) = \sum_{l=0}^{\infty} c_l P_l(x) \quad (45)$$

$$(46)$$

where

$$c_m = \frac{2m+1}{2} \int_{-1}^1 f(x)P_l(x)dx \quad (47)$$

9. Associated Legendre functions:

$$P_l^m(x) = (1-x^2)^{m/2} \frac{d^m}{dx^m} P_l(x) \quad (48)$$

satisfy the equation

$$(1-x^2)y'' - 2xy' + \left[l(l+1) - \frac{m^2}{1-x^2} \right] y = 0 \quad (49)$$

For *each* m , they a set of orthogonal functions on $(-1, 1)$, with normalization:

$$\int_{-1}^1 [P_l^m(x)]^2 dx = \frac{2}{2l+1} \frac{(l+m)!}{(l-m)!} \quad (50)$$

10. Bessel equation

$$x^2 y'' + xy' + (x^2 - p^2) y = 0 \quad (51)$$

has solutions (Bessel functions):

$$J_p(x) = \sum_{n=0}^{\infty} \frac{(-1)^n}{\Gamma(n+1)\Gamma(n+1+p)} \left(\frac{x}{2}\right)^{2n+p} \quad (52)$$

and

$$N_p(x) = \frac{\cos(\pi p)J_p(x) - J_{-p}(x)}{\sin \pi p} \quad (53)$$

11. Asymptotic values:

$$J_0(0) = 1 \quad (54)$$

$$J_{n \neq 0}(0) = 0 \quad (55)$$

$$J_{n=0,1,2\dots}(\infty) = 0 \quad (56)$$

12. Recursion relations for Bessel functions

$$\frac{d}{dx} [x^p J_p(x)] = x^p J_{p-1}(x) \quad (57)$$

$$\frac{d}{dx} [x^{-p} J_p(x)] = -x^{-p} J_{p+1}(x) \quad (58)$$

$$J_{p-1}(x) + J_{p+1}(x) = \frac{2p}{x} J_p(x) \quad (59)$$

$$J_{p-1}(x) - J_{p+1}(x) = 2J'_p(x) \quad (60)$$

$$J'_p(x) = -\frac{p}{x} J_p(x) + J_{p-1}(x) = \frac{p}{x} J_p(x) - J_{p+1}(x) \quad (61)$$

13. Other equations with Bessel function solutions

$$y'' + \frac{1-2a}{x} y' + \left[(bcx^{c-1})^2 + \frac{a^2 - p^2 c^2}{x^2} \right] y = 0 \quad (62)$$

has the solution

$$y = x^a Z_p(bx^c), \text{ where } Z = J, N \quad (63)$$

and

$$y = J_p(Kx) \text{ and } N_p(Kx) \quad (64)$$

satisfy the equation

$$x(xy')' + (K^2 x^2 - p^2) y = 0 \quad (65)$$

14. Orthogonality of Bessel functions:

$$\int_0^1 x J_p(\alpha x) J_p(\beta x) = \begin{cases} 0 & \text{if } \alpha \neq \beta, \\ \frac{1}{2} J_{p+1}^2(\alpha) = \frac{1}{2} J_{p-1}^2(\alpha) = \frac{1}{2} J_p'^2(\alpha) & \text{if } \alpha = \beta \end{cases} \quad (66)$$

where α and β are zeroes of $J_p(x)$.

Chapter 13 of Boas (Partial Differential Equations)

1. Laplace equation in two-dimensional rectangular/Cartesian coordinates (for example, for steady-state temperature):

$$\frac{\partial^2}{\partial x^2} T(x, y) + \frac{\partial^2}{\partial y^2} T(x, y) = 0 \quad (67)$$

has basis functions (i.e., general solution is a suitable *combination* of these):

$$T(x, y) = \left\{ \begin{array}{l} e^{kx} \\ e^{-kx} \end{array} \right\} \left\{ \begin{array}{l} \sin ky \\ \cos ky \end{array} \right\} \text{ or } \left\{ \begin{array}{l} \sin kx \\ \cos kx \end{array} \right\} \left\{ \begin{array}{l} e^{ky} \\ e^{-ky} \end{array} \right\} \quad (68)$$

2. Diffusion equation in one dimension

$$\frac{\partial^2 u}{\partial x^2} = \frac{1}{\alpha^2} \frac{\partial u}{\partial t} \quad (69)$$

has basis functions

$$u = e^{-k^2 \alpha^2 t} \left\{ \begin{array}{l} \sin kx \\ \cos kx \end{array} \right\} \quad (70)$$

3. Schroedinger equation in one dimension for a *free* particle (i.e., no potential):

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} = i\hbar \frac{\partial \Psi}{\partial t} \quad (71)$$

has basis functions

$$\Psi = \left\{ \begin{array}{l} \sin kx \\ \cos kx \end{array} \right\} e^{-iEt/\hbar} \quad (72)$$

4. Wave equation in circular coordinates:

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial z}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 z}{\partial \theta^2} = \frac{1}{v^2} \frac{\partial^2 z}{\partial t^2} \quad (73)$$

has basis functions:

$$z = \left\{ \begin{array}{l} J_n(Kr) \\ N_n(Kr) \end{array} \right\} \left\{ \begin{array}{l} \sin n\theta \\ \cos n\theta \end{array} \right\} \left\{ \begin{array}{l} \sin Kvt \\ \cos Kvt \end{array} \right\} \quad (74)$$

5. Laplace equation in spherical coordinates

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial u}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial u}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 u}{\partial \phi^2} = 0 \quad (75)$$

has basis functions (where l is a non-negative integer, with $-l \leq m \leq +l$)

$$u = \left\{ \begin{matrix} r^l \\ r^{-l-1} \end{matrix} \right\} P_l^m(\cos \theta) \left\{ \begin{matrix} \sin m\phi \\ \cos m\phi \end{matrix} \right\} \quad (76)$$

Chapter 14 of Boas (Functions of a Complex Variable)

1. Basics of complex-valued functions of complex variable

$$f(z) = f(x + iy) = u(x, y) + iv(x, y) \quad (77)$$

$$f'(z) = \frac{df}{dz} = \lim_{\Delta z \rightarrow 0} \frac{\Delta f}{\Delta z} \quad (78)$$

2. If $f(z)$ is analytic in a region (i.e., has a unique derivative at every point), then

$$\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}, \quad (79)$$

$$\frac{\partial v}{\partial x} = -\frac{\partial u}{\partial y} \quad (80)$$

(Cauchy-Reimman conditions) and it's converse: if $u(x, y)$ and $v(x, y)$ satisfy these conditions, then $f(z) = u + iv$ is analytic.

3. If $f(z)$ is analytic in a region R , then it has derivatives of all orders at points inside R and thus it can be expanded in a Taylor series about any point z_0 in R . This power series converges *inside* circle C about z_0 that extends to the nearest singularity point (i.e., C *just* touches the boundary of R).

4. If $f(z) = u + iv$ is analytic in a region, then u and v satisfy (two-dimensional) Laplace's equation in the region. And, conversely, *any* function u (or v) satisfying Laplace's equation is the real (or imaginary) part of an analytic function $f(z)$.

5. Cauchy's theorem: if $f(z)$ is analytic inside and on a closed curve C , then

$$\int_C f(z) dz = 0, \text{ around } C \quad (81)$$

6. Cauchy's integral formula: if $f(z)$ is analytic inside and on a closed curve C , then

$$f(a) = \frac{1}{2\pi i} \int_C \frac{f(z)}{z - a} dz, \text{ around } C \quad (82)$$

where $z = a$ is a point *inside* C .

7. Laurent series: Let C_1 and C_2 be two circles with center at z_0 . If $f(z)$ is an analytic function in the region R between $C_{1,2}$, then it can be expanded in a convergent series in R

$$f(z) = a_0 + a_1(z - z_0) + a_2(z - z_0)^2 + \dots + \frac{b_1}{z - z_0} + \frac{b_2}{(z - z_0)^2} + \dots \quad (83)$$

associated with which are the following definitions:

- (i) If all the b 's are zero, then $f(z)$ is analytic at $z = z_0$ (regular point);
- (ii) If $b_n \neq 0$, but all the subsequent b 's are zero, then $f(z)$ is said to have a *pole* of order n at $z = z_0$. If $n = 1$ here, then $f(z)$ has a simple pole at $z = z_0$;
- (iii) If there are infinite number of b 's which are different than zero, then $f(z)$ has an *essential* singularity at $z = z_0$;
- (iv) The coefficient b_1 of $1/(z - z_0)$ is called the *residue* of $f(z)$ at $z = z_0$.

8. Residue theorem:

$$\int f(z)dz \text{ (around } C) = 2\pi i \cdot (\text{sum of residues of } f(z) \text{ inside } C) \quad (84)$$

where we go *counter*-clockwise around C .

9. Methods of finding residues of $f(z)$:

(A) coefficient b_1 in Laurent series about $z = z_0$;

(B) Simple pole:

$$R(z_0) = \lim_{z \rightarrow z_0} (z - z_0) f(z) \quad (85)$$

and if $f(z) = g(z)/h(z)$, then

$$R(z_0) = \frac{g(z_0)}{h'(z_0)} \text{ if } \begin{cases} \text{if } g(z_0) = \text{finite const.} \\ h(z_0) = 0, h'(z_0) \neq 0 \end{cases} \quad (86)$$

(C) Multiple poles: multiply $f(z)$ by $(z - z_0)^m$, where m is an integer $\geq n$ (order of pole), differentiate the result $(m - 1)$ times, divide by $(m - 1)!$, and evaluate the resulting expression at $z = z_0$.

10. Definite integrals using residue theorem:

(i) Change of variables;

(ii) If $P(x)$ and $Q(x)$ are polynomials with degree of $Q \geq$ degree of $P + 2$ and if Q has no real zeroes, then

$$\int_{-\infty}^{\infty} \frac{P(x)}{Q(x)} = 2\pi i. \left(\text{sum of residues of } \frac{P(z)}{Q(z)} \text{ in upper half-plane} \right) \quad (87)$$

(iii) If $P(x)$ and $Q(x)$ are polynomials with degree of $Q \geq$ degree of $P + 1$ and if Q has no real zeroes, then

$$\int_{-\infty}^{\infty} \frac{P(x)}{Q(x)} e^{imx} = 2\pi i. \left(\text{sum of residues of } \frac{P(z)}{Q(z)} e^{imz} \text{ in upper half-plane} \right) \quad (88)$$

where $m > 0$.

(iv) Poles on boundary:

$$\int f(z) dz \text{ (around } C) = 2\pi i. \left(\text{sum of residues at simple poles inside } C + \frac{1}{2} \text{ sum of residues of poles on the boundary} \right) \quad (89)$$

(v) Branch cuts: for integrals involving *fractional* powers (or logarithm) of x (and thus z), we have to choose contour such that we stay on one branch of the fractional power (say, angle of z between 0 and 2π) so that the function is *single*-valued.

(vi) Argument principle:

$$N - P = \frac{1}{2\pi i} \int \frac{f'(z)}{f(z)} dz \text{ (around } C) = \frac{1}{2\pi} \Theta_C \quad (90)$$

where N and P are the number of zeroes and poles, respectively, of $f(z)$ inside C and Θ_C is the change in angle of $f(z)$ around C .

11. Nature of $f(Z)$ at $Z = \infty$: it is a pole of order 2 if $f(1/z)$ is the same at $z = 0$ etc.

12. Residue at infinity:

$$\left(\text{residue of } f(Z) \text{ at } Z = \infty \right) = - \left(\text{residue of } \frac{1}{z^2} f \left(\frac{1}{z} \right) \text{ at } z = 0 \right) \quad (91)$$